

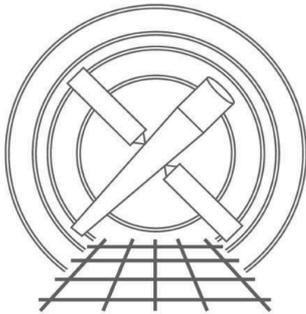
Detecting Diffractive Lensing in Astrophysical Gravitational Waves

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With Shun-Sheng Li, Barak Zackay,
Shude Mao and Youjun Lu

Einstein Fellowship Symposium @ Harvard CfA

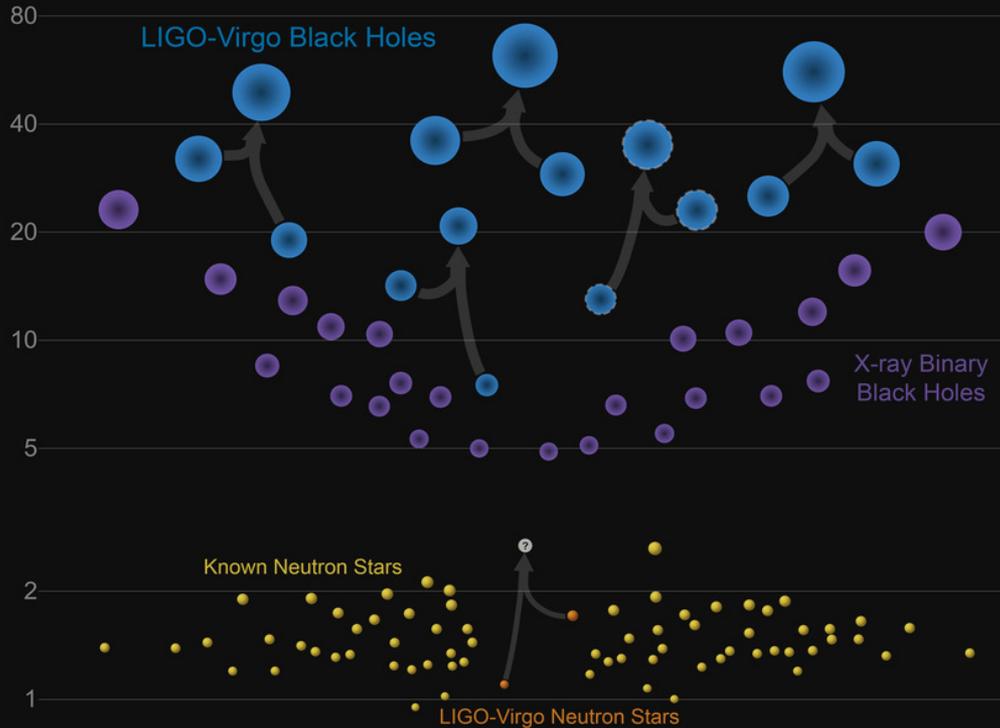
Oct 2018



Gravitational Waves: A New Window Into the Universe

Masses in the Stellar Graveyard

in Solar Masses



LIGO Hanford/Livingston



Virgo

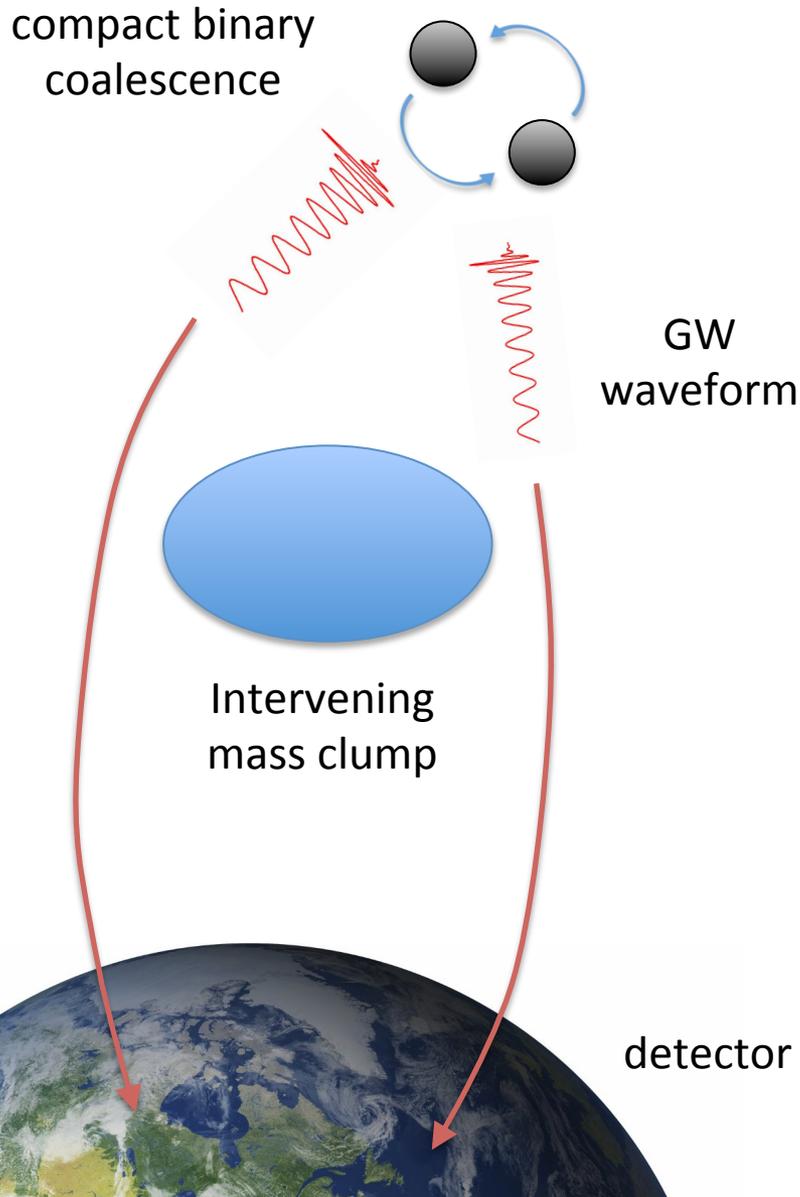
KAGRA



LIGO-India

Space-based observatory (LISA)
Third generation detector

Gravitational lensing of gravitational waves



Wang, Stebbins & Turner 96'
Li, Mao, Zhao & Lu 18'

- ◆ Regime of geometrical lensing

$$\lambda_{GW} \ll R_{de}$$

- ◆ Amplification of strain

$$h' = \sqrt{\mu} h$$

- ◆ Wave frequency unchanged

Observational difficulty with geometrical lensing

Dai, Venumadhav & Sigurdson 2017

Ng, Wang, Broadhurst & Li 2017

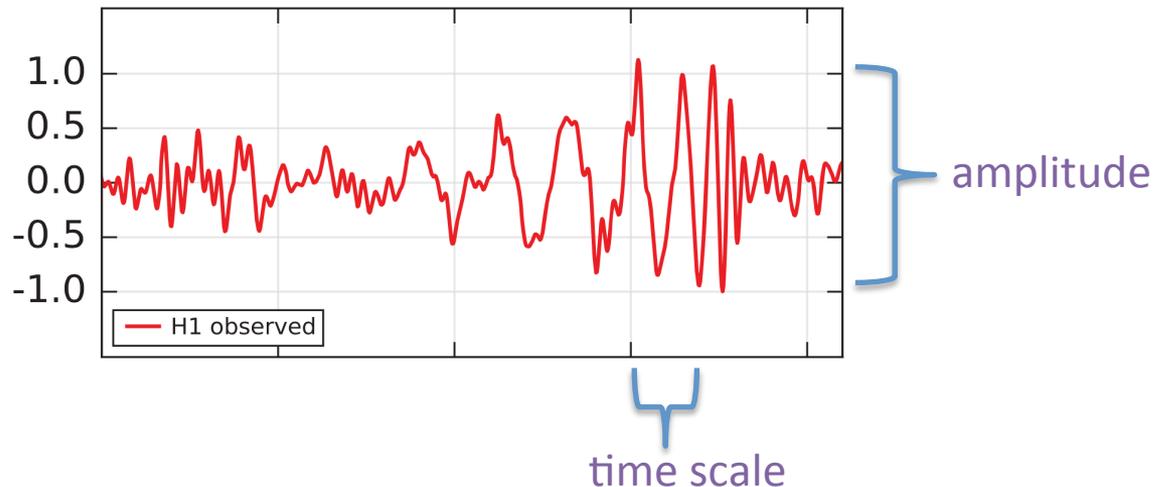
Broadhurst, Diego & Smoot 2018

Oguri 2018

Apparent mass scale M'

Apparent source redshift z'

$$M(1+z) = M'(1+z') \quad \frac{\sqrt{\mu}}{d_L(z)} = \frac{1}{d_L(z')}$$



Without EM observations, magnification cannot be recovered from a single lensed image

Topological (Morse) phase shift for **flipped** images; however degenerate with orbital orientation; Dai & Venumadhav 2017

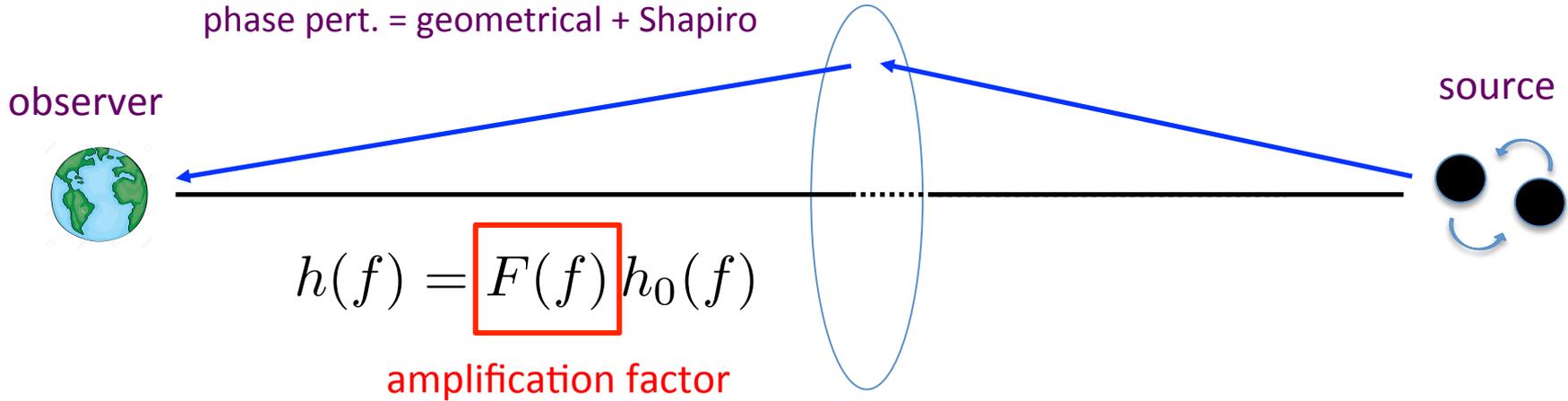
Multiple lensed images resolvable in time domain

Require fine-tuned impact parameters and large column density

Can be produced by cluster/galaxy lenses; difficult for low-mass lenses

Lensing in wave-diffraction regime

Takahashi & Nakamura 2003
Takahashi PhD thesis



Limit of geometric optics

$$w = 2\pi f (1 + z_L) G M_L / c^2 \gg 1$$

(if only one image) $F_{\text{geo}}(f) \longrightarrow \sqrt{\mu_{\text{geo}}} e^{i 2\pi f (1+z_L) \tau_{\text{geo}}}$ time delay

Regime of wave diffraction

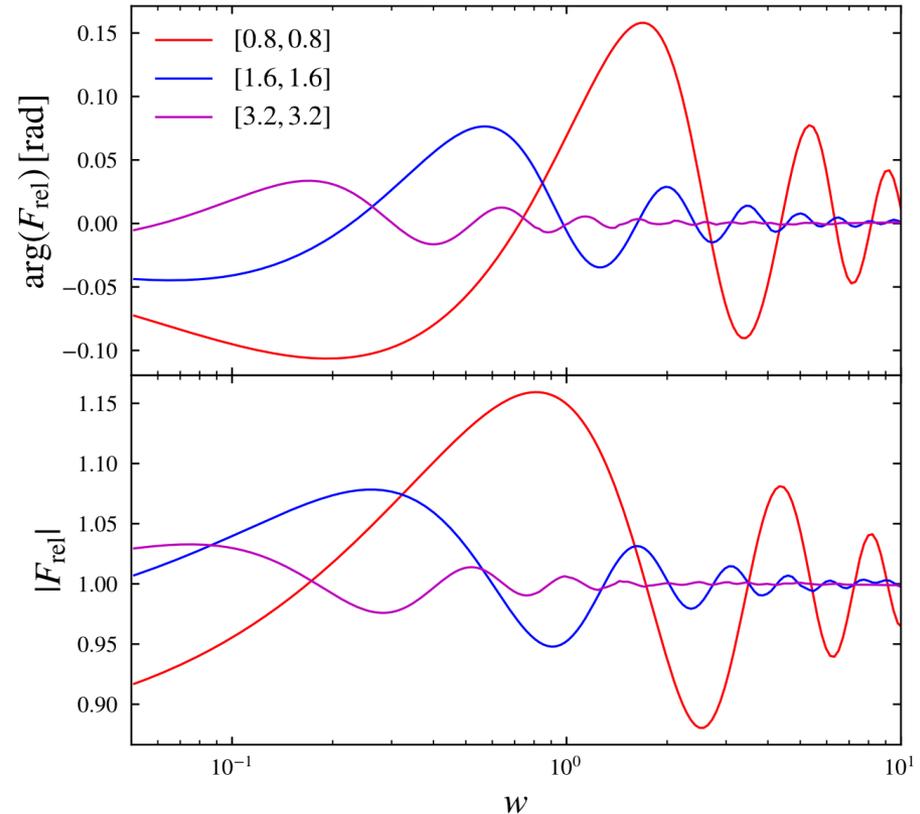
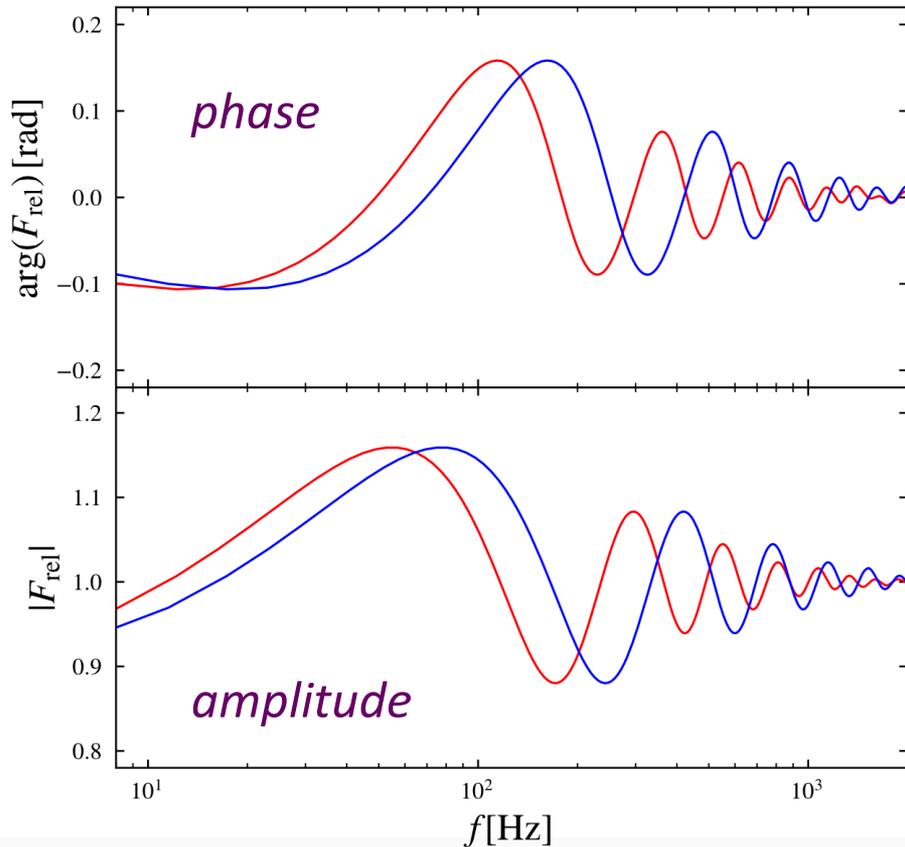
(frequency-independent)
amplification

$$w = 2\pi f (1 + z_L) G M_L / c^2 \sim \mathcal{O}(1)$$

Define $F_{\text{rel}}(f) := F(f) / F_{\text{geo}}(f)$ Encode information of waveform distortion !

Amplitude and Phase Modulations

Ground-based band $f \sim 10\text{-}1000$ Hz sensitive to (interestingly) small lens masses $M \sim 100\text{-}1000 M_{\text{sun}}$ Dai, Li, Zackay, Mao & Lu 2018



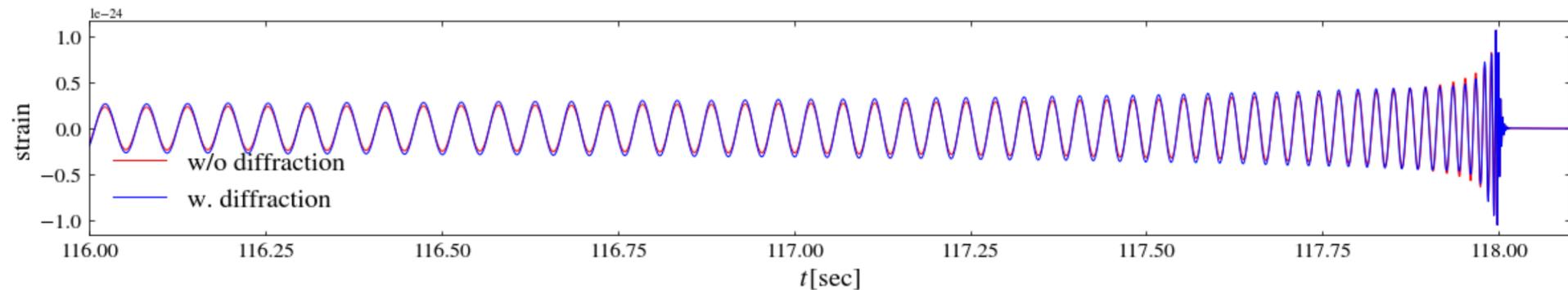
Pseudo-Jaffe lens $\sigma_v \sim 1 - 2$ km/s
 Impact parameter $\theta_E \sim 4\pi (\sigma_v/c)^2$

Size of modulation inversely proportional to the impact parameter
 Lens mass scale increases

Diffraction signature is subtle

Small modulus and phase perturbations $\sim 10\text{-}20\%$ or even smaller!

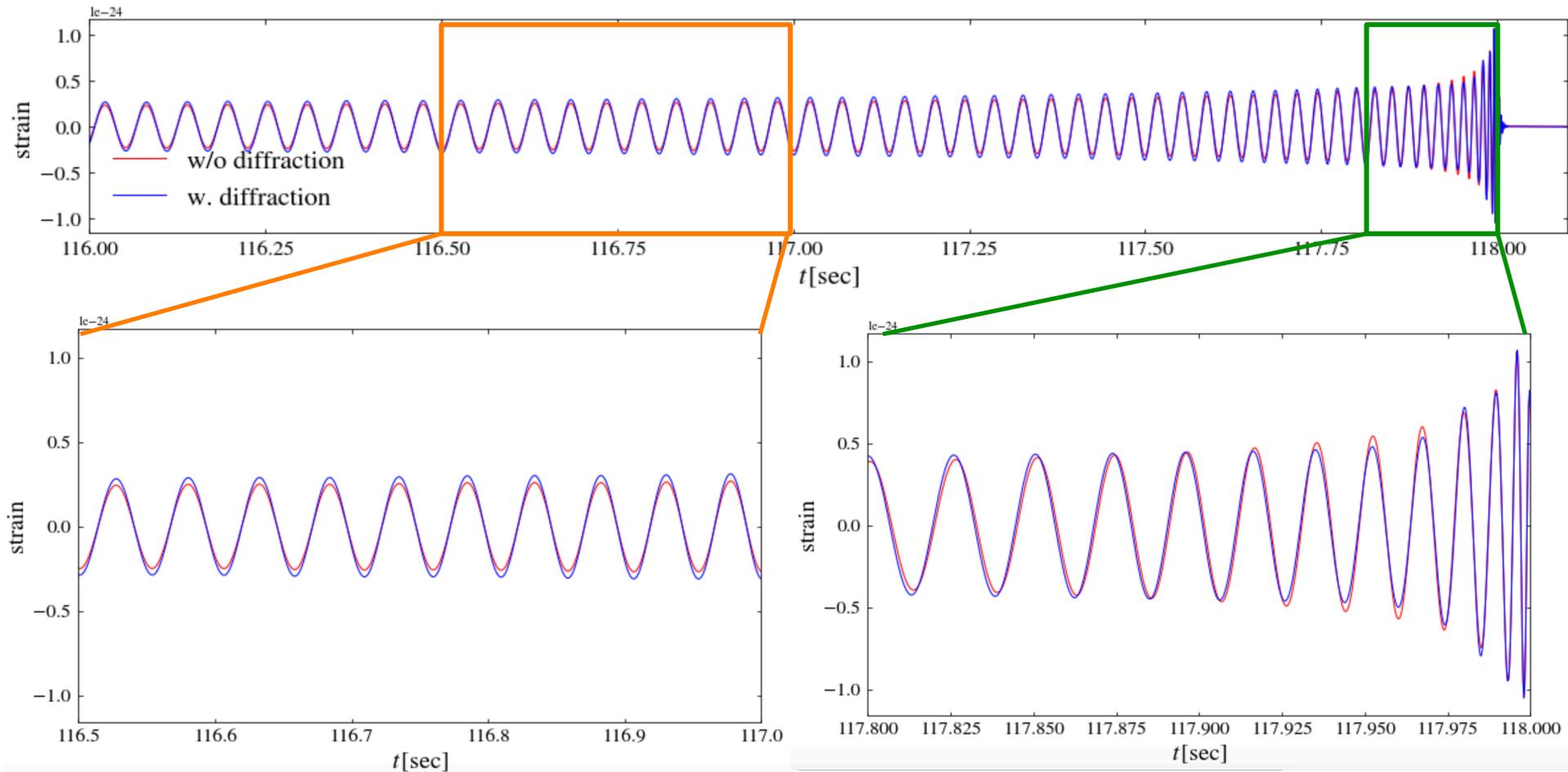
Can you see the amplitude/phase modulations?



Diffraction signature is subtle

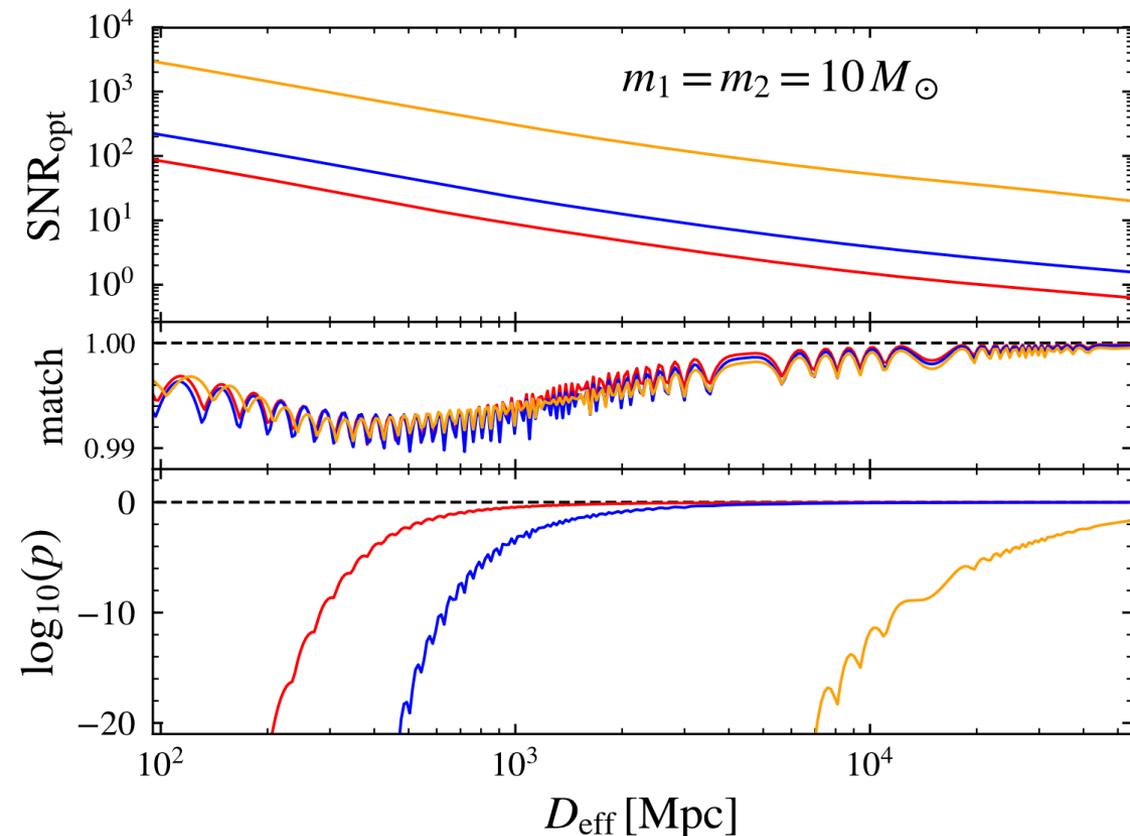
Small modulus and phase perturbations $\sim 10\text{-}20\%$ or even smaller!

Can you see the amplitude/phase modulations? Hmm ... not so impressive ...



Match with unlensed templates is (nearly) unaffected.

Diffraction signature still detectable through **the improvement in the likelihood** when **amplitude/phase modulations** are included into the waveform.



aLIGO_MID_LOW

aLIGO_DESIGN

Einstein Telescope

$$\text{match} := \frac{\langle h_L | h_{\text{BF}} \rangle}{\sqrt{\langle h_L | h_L \rangle \langle h_{\text{BF}} | h_{\text{BF}} \rangle}}$$

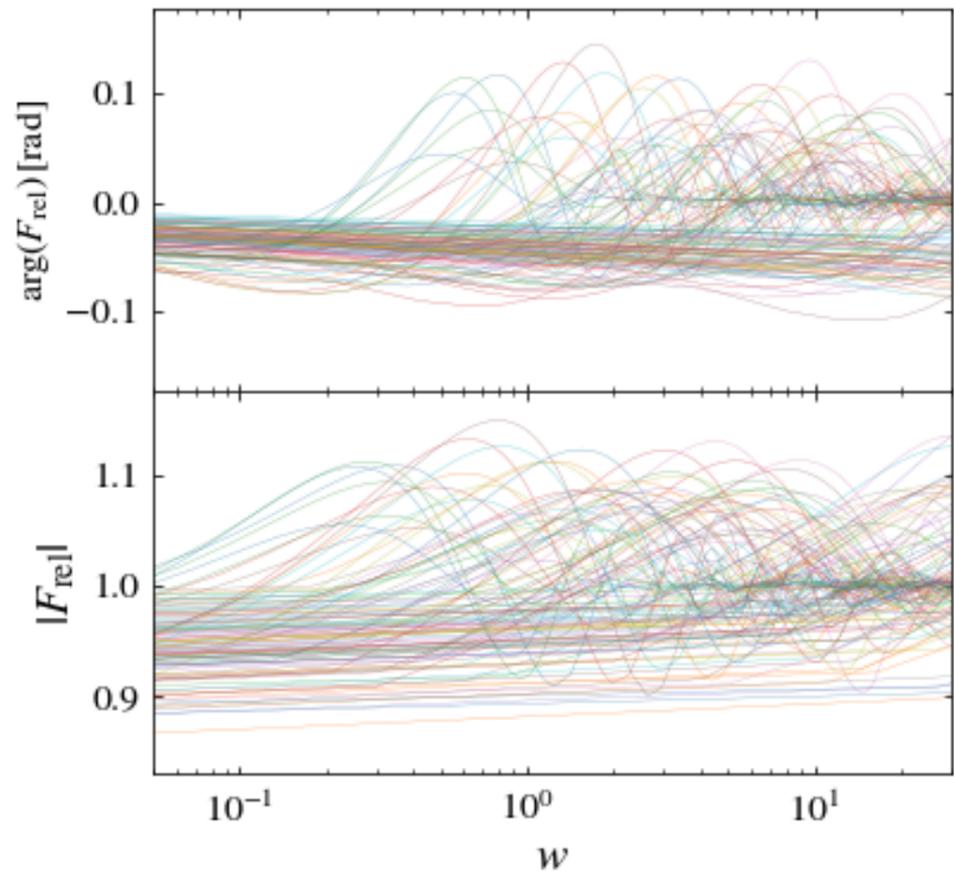
$$\begin{aligned} \ln p &= -(\text{SNR}_{\text{opt}}^2 - \text{SNR}_{\text{unlen}}^2) / 2 \\ &= -\frac{1}{2} \left(\langle h_L | h_L \rangle - \frac{\langle h_L | h_{\text{BF}} \rangle^2}{\langle h_{\text{BF}} | h_{\text{BF}} \rangle} \right) \end{aligned}$$

Matched filtering and some practical difficulties

- ◆ Matched filtering requires the precise knowledge of $F(f)$

e.g. Takahashi & Nakamura 2003,
Cao+ 2014, Jung & Shin 2017

- ◆ $F(f)$ depends on too many parameters: **lens profile, distances, impact parameter, etc.**
- ◆ The correct lens profile to use is unknown.
- ◆ Have to search with a large number of templates. **Look-elsewhere effect** needs to be quantified.

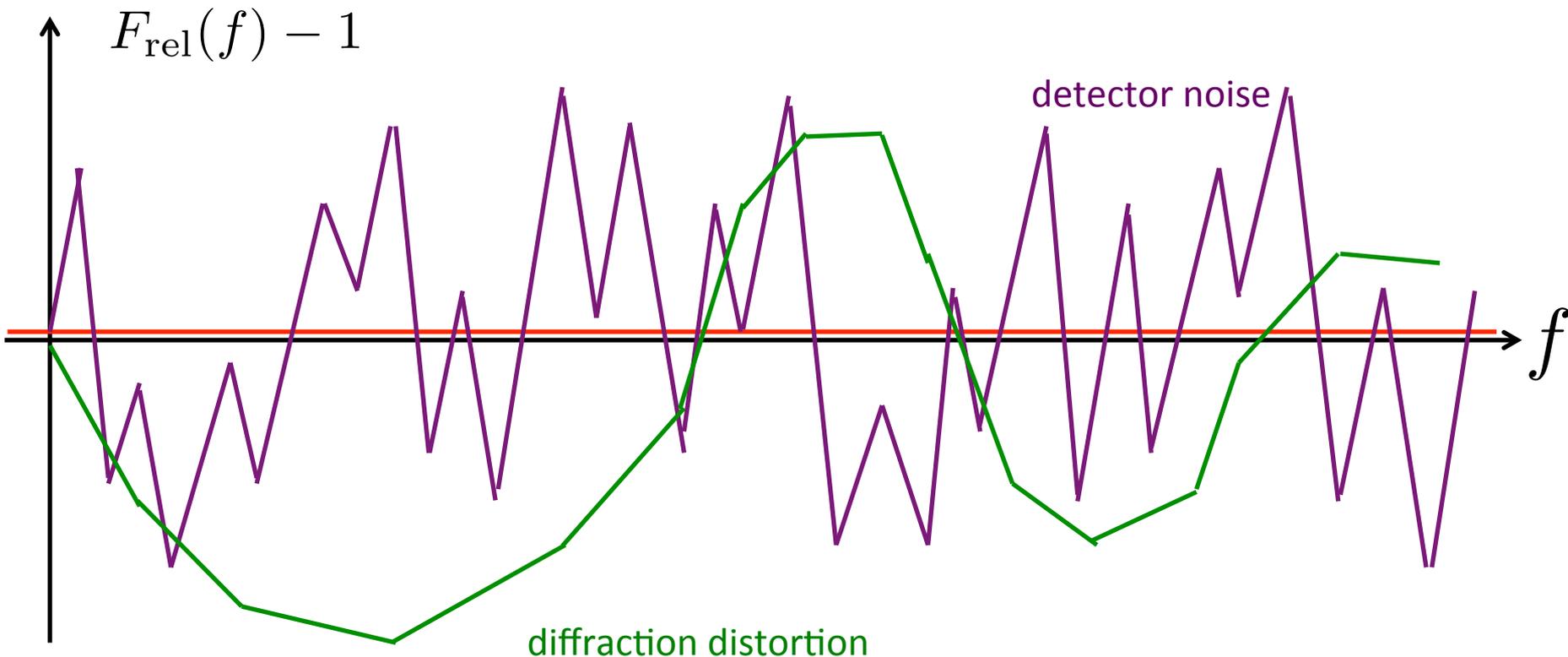


An agnostic method based on dynamic programming

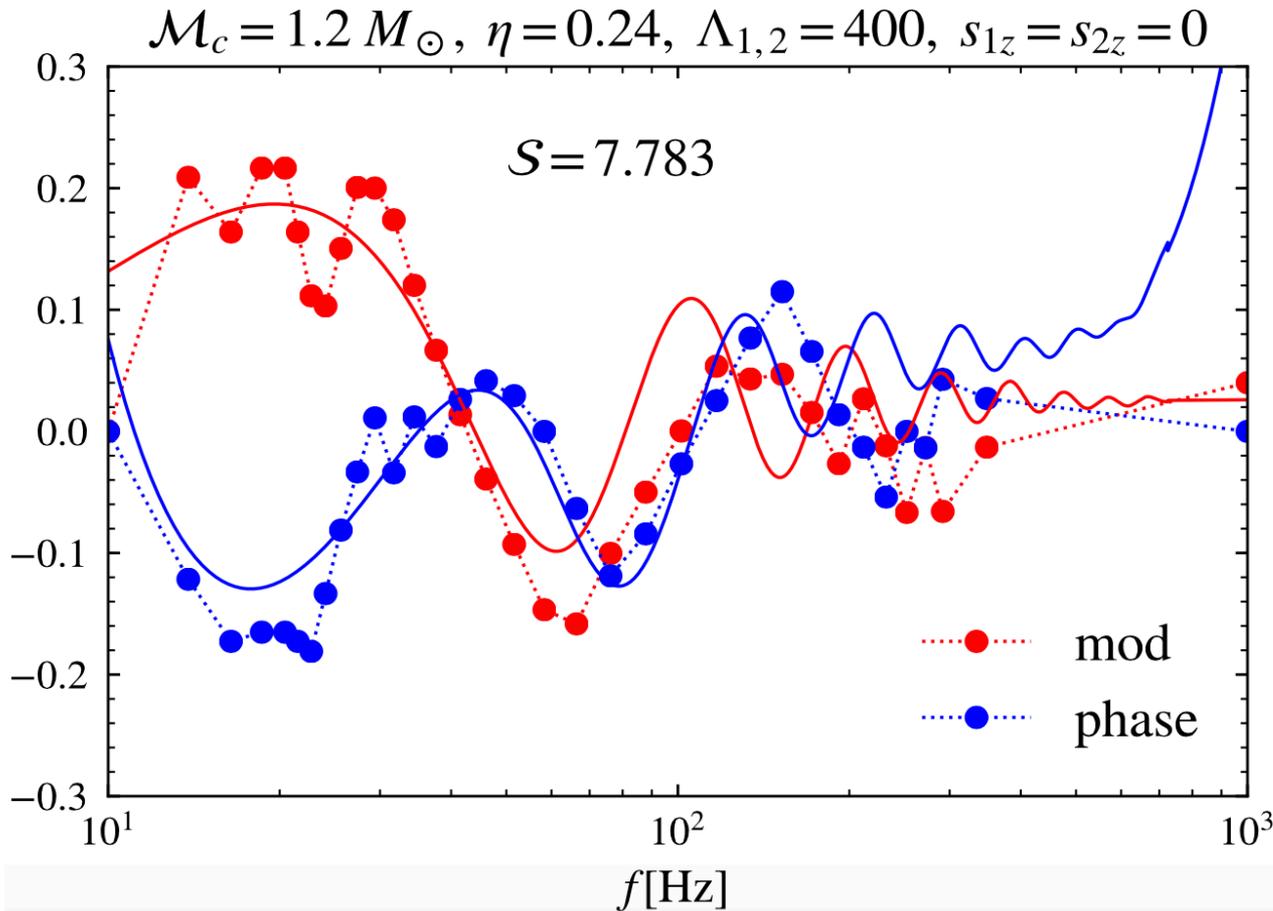
Dai+ 2018

$$\mathcal{S} := \int \mathcal{D}g(f) \underbrace{\mathcal{P}[g(f)]}_{\substack{\text{prior} \\ \text{(Markovian)}}} \prod_{a=1}^{N_d} \underbrace{\frac{P[s_a(f)|g(f) h_{\text{BF},a}(f)]}{P[s_a(f)|h_{\text{BF},a}(f)]}}_{\text{likelihood improvement}}$$

$h_{\text{BF}}(f)$ is the **best-fit unlensed** waveform



Reconstructing the modulations



Pseudo-Jaffe lens
at $z = 0.1$

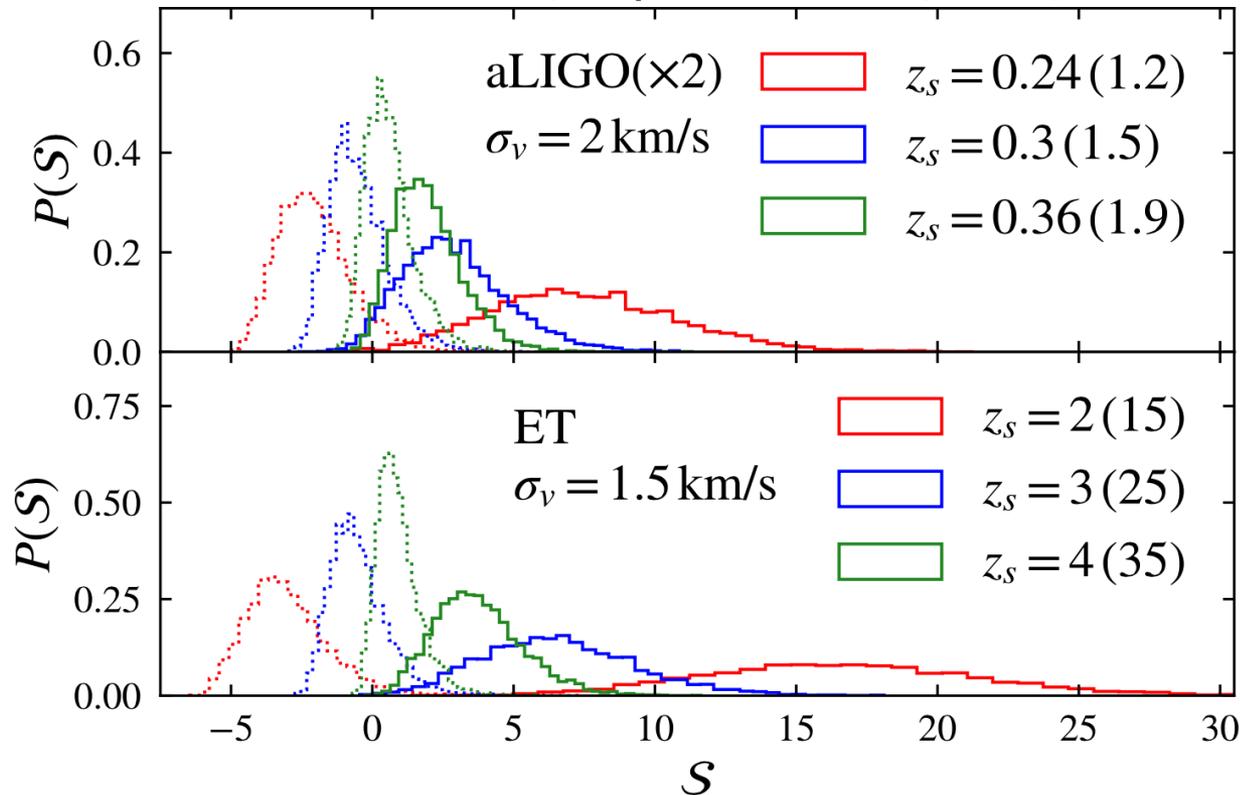
Velocity dispersion
 $\sigma_v = 2$ km/s

NS-NS merger at
 $z = 0.2$

Can only detect the part of the modulation signal that is not degenerate with source parameters !

Observational prospect: BH mergers are promising!

$$\mathcal{M}_c = 30 M_\odot, \eta = 0.24, s_{1z} = s_{2z} = 0$$



Assume pseudo-Jaffe halos; mass enclosed within the Einstein radius ~ 100 --- $1000 M_{\text{sun}}$

aLIGO Hanford/Livingston can probe out to $z \sim 0.2 - 0.3$.

Further improves **after more detectors join** (Virgo, KAGRA, LIGO-India, etc)

3rd generation detector will be very powerful: $z \sim 2 - 4$

Discussion

- ◆ Science case: test **CDM theory** on sub-galactic scales ?!
 - Probe inner region of $M \sim 10^4 - 10^6 M_{\text{sun}}$ DM halos
 - 3rd gen. detector can use BBHs out to $z \sim 2-4$
 - Assume nearly log-flat halo mass function; lensing optical depth $\sim \text{few} * 10^{-3}$ if $r_E \sim 1 \text{ pc}$
 - Enough enclosed mass? Small halos show steeper inner profiles than NFW.
e.g. [Dutton & Maccio 2014](#)
 - **Galaxy lensing events** particularly interesting to look at

Thank you!

◆ Degeneracy with **spin-precessing** or **eccentricity** effects?

- Precession can induce amplitude/phase modulations in the frequency-domain waveform. e.g. [Apostolatos, Cutler, Sussman & Thorne, 1994](#); [Klein, Cornish & Yunes 2013](#)
- Modulation frequency (**~ tens of precession cycles in band**) is typically higher than diffraction; amplitude modulation more significant than phase modulation.
- Detailed study would be very valuable; need accurate waveforms.